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Optimal multigrid convergence by elliptic/hyperbolic splitting

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Abstract

We describe a multigrid method for solving the steady Euler equations that is optimal in the sense of requiring $O(N)$ operations till convergence, where N is the number of unknowns. The method relies on an elliptic/hyperbolic decomposition achieved by local preconditioning. The splitting allows the embedded advection equations to be treated with streamwise semicoarsening rather than full coarsening, which would not be effective. A simple 2-D numerical computation is presented as proof of concept. A convergence study indicates the split method has complexity $O(N)$ over a wide range of grid spacings and Mach numbers, while the use of full coarsening for all equations makes the complexity deteriorate to almost $O(N^{1.5})$.

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1. Introduction

We describe a multigrid method for solving the steady Euler equations that is optimal in the sense of requiring $O(N)$ operations till convergence, where N is the number of unknowns. The method relies on an elliptic/hyperbolic decomposition achieved by local preconditioning. This method means a significant simplification and computational savings compared to the method of Darmofal and Siu [2], who also used local preconditioning in order to reach $O(N)$ complexity. Their multigrid method required semicoarsening in all directions, which is computationally expensive, especially for three-dimensional problems. Isotropic semicoarsening, though robust, was understood by those and the current authors to be “overkill” [16]: *for optimal convergence, semicoarsening need only be applied to the embedded advection equations, and only in the flow direction.* Such “streamwise” semicoarsening requires the splitting of the residual into advective and acoustic components, which happens to be a special benefit of the van Leer–Lee–Roe [17] local preconditioning. Below we use local preconditioning for removing local stiffness as well as for equation decomposition, and we compare the rate of convergence to subsonic flow solutions for the new “elliptic/hyperbolic multigrid” and standard (full-coarsening) multigrid methods.

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2. Toward $O(N)$ complexity

In a long sequence of papers running from 1989 [15] to 1999 [2] van Leer and collaborators, later joined by Darmofal and collaborators, demonstrated that it is possible to achieve convergence to steady Euler solutions in $O(N)$ operations using an explicit single-grid scheme augmented by multigrid relaxation, provided:

- (1) local preconditioning is used to remove the local stiffness caused by the spread of the characteristic speeds;
- (2) the single-grid scheme is designed to efficiently remove all removable high-frequency errors (some errors are not removable when the flow is grid-aligned);
- (3) a multigrid method more powerful than standard full-coarsening multigrid is used to overcome the alignment problem.

An overview of this research effort has been given by van Leer and Darmofal [16]; in this paper we shall only present supplemental information needed to understand the current method. This includes a discussion and comparison of the benefits of isotropic and streamwise semicoarsening, and a discussion of elliptic/hyperbolic splitting.

In this paper we restrict ourselves to the subsonic 2-D Euler equations and their numerical solution. (The supersonic Euler equations are simpler in that they are fully hyperbolic and can be handled by one semicoarsening strategy.) There is a complication when going to 3-D Euler, caused by the inability to separate the advection equation for streamwise vorticity from the acoustic equations by algebraic preconditioning [12,17]; we will return to this subject in the last section.

3. Baseline Euler solver

Our aim is the efficient computation of steady-state solutions to the Euler equations in two dimensions:

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} + \frac{\partial \mathbf{G}}{\partial y} = \mathbf{0}, \quad (1)$$

$$\mathbf{U} = \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{bmatrix}, \quad \mathbf{F} = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ \rho uH \end{bmatrix}, \quad \mathbf{G} = \begin{bmatrix} \rho v \\ \rho uv \\ \rho v^2 + p \\ \rho vH \end{bmatrix}, \quad (2)$$

where ρ is the density, u and v are the velocity components in the x - and y -direction, respectively, and p is the static pressure. The specific total energy and enthalpy are given by

$$E = \frac{1}{\gamma - 1} \frac{p}{\rho} + \frac{1}{2} (u^2 + v^2), \quad (3)$$

$$H = \frac{\gamma}{\gamma - 1} \frac{p}{\rho} + \frac{1}{2} (u^2 + v^2). \quad (4)$$

We discretize this system with a cell-centered finite-volume algorithm on a structured quadrilateral grid. The cell-residual \mathbf{Res} is defined as a numerical approximation of the integral over the cell of the spatial operator, and can be written as

$$\mathbf{Res} = - \sum_{k=1}^4 \mathbf{\Phi}_k \Delta s_k, \quad (5)$$

where Δs_k is the length of cell face k and $\mathbf{\Phi}_k$ is the outward numerical flux normal to the face. In this work, we use Roe's approximate Riemann solver, with a modification required for stability in the presence of preconditioning, and for preserving solution accuracy at low Mach numbers

$$\mathbf{\Phi}_k = \frac{1}{2} [\mathbf{\Phi}(\mathbf{U}_L) + \mathbf{\Phi}(\mathbf{U}_R)] - \frac{1}{2} \hat{\mathbf{P}}^{-1} |\hat{\mathbf{P}} \hat{\mathbf{A}}| (\mathbf{U}_R - \mathbf{U}_L). \quad (6)$$

Here $\hat{\mathbf{P}}$ is a local preconditioning matrix and $\hat{\mathbf{A}}$ is the flux Jacobian, each evaluated at the cell face using Roe's average state. To achieve second-order accuracy, Van Leer's κ -scheme [14] is employed with $\kappa = 0$ and without limiting. We purposely omit the limiting because this highly nonlinear technique is known to stall convergence [1] and might impede our proof of concept. Without limiting, solutions would not be monotonic near shocks for this choice of κ ; to the subsonic test problem used below this does not apply.

The residual defined as above drives the evolution of the cell-average $\bar{\mathbf{U}}$ of \mathbf{U} according to the preconditioned equation

$$\frac{\partial \bar{\mathbf{U}}}{\partial t} = \frac{1}{S} \mathbf{P}(\bar{\mathbf{U}}) \mathbf{Res}, \quad (7)$$

where S is the area of the cell.

To reach a steady state we integrate this system in time using a multistage time-stepping optimized for high-frequency damping, as required for use in multigrid relaxation. The particular scheme employed here is the four-stage optimized smoothing scheme developed for $\kappa = 0$ by Lynn and van Leer [4,5]; it has the form

$$\begin{aligned} \bar{\mathbf{U}}^{(1)} &= \bar{\mathbf{U}}^n + \frac{\beta^{(1)} \Delta t}{S} \mathbf{P}^n \mathbf{Res}^n, \\ \bar{\mathbf{U}}^{(2)} &= \bar{\mathbf{U}}^n + \frac{\beta^{(2)} \Delta t}{S} \mathbf{P}^{(1)} \mathbf{Res}^{(1)}, \\ \bar{\mathbf{U}}^{(3)} &= \bar{\mathbf{U}}^n + \frac{\beta^{(3)} \Delta t}{S} \mathbf{P}^{(2)} \mathbf{Res}^{(2)}, \\ \bar{\mathbf{U}}^{n+1} &= \bar{\mathbf{U}}^n + \frac{\beta^{(4)} \Delta t}{S} \mathbf{P}^{(3)} \mathbf{Res}^{(3)}. \end{aligned} \quad (8)$$

We have

$$\Delta t = 2\nu S / \sum_{k=1}^4 |q_k| ds_k. \quad (9)$$

Here q_k is the flow-velocity component normal to cell face k (the preconditioning makes all characteristic speeds equal to the flow velocity), the CFL-number ν has the value 1.4008, and $\beta^{(1)} = 0.1299$, $\beta^{(2)} = 0.2940$, $\beta^{(3)} = 0.5604$, $\beta^{(4)} = 1.0000$.

4. Semicoarsening

Multigrid relaxation on a collection of semicoarsened grids was developed by Mulder [9] for increased robustness, in particular, when the flow is aligned over appreciable portions of the grid. In that case it is neither possible nor desirable for a single-grid scheme to damp error modes with high frequencies in the

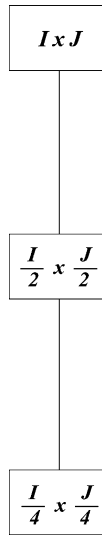


Fig. 1. Full coarsening.

stream-normal direction for advected quantities such as entropy. The standard full-coarsening multigrid process (see Fig. 1) then suffers because streamwise low-frequency errors coupled to those modes get averaged out during full coarsening and are not addressed at all. This slows down overall convergence and makes $O(N)$ complexity elusive. Streamwise semicoarsening will make the multigrid process effective in removing streamwise low-frequency error modes from the advected quantities. Since it may not be known in advance if, where, and how the flow aligns with the grid, Mulder reasoned that a robust, general-purpose strategy would be to use contributions from a matrix of grids semicoarsened in all possible directions (see Fig. 2). This will catch all high-frequency advective errors, while also being a powerful strategy for removing acoustic error modes. Thus, there is no need to decompose the residual into acoustic (elliptic) and

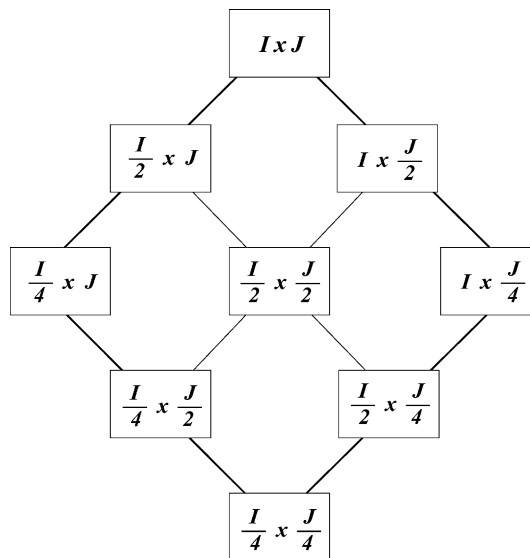


Fig. 2. Mulder's isotropic semi-coarsening.

advective (hyperbolic) components. Furthermore, because of the powerful one-dimensional mode-removal in the corrections on semi-coarsened grids, the demands on the single-grid scheme regarding smoothing are significantly reduced: only error modes with high frequencies in *all* directions need to be damped on each grid. Multistage marching schemes can be designed to do so very effectively, provided the residuals are first treated with local preconditioning. This clusters the eigenvalues of the spatial operator and makes it possible to eliminate the corresponding error modes regardless of the Mach number.

Although powerful, such isotropic semicoarsening is computationally too expensive. Compared to single-grid relaxation, the isotropic semicoarsening adds a factor 4 to the work in 2D and 8 in 3D which are significantly larger than the corresponding factors 4/3 and 8/7 for full coarsening. On the other hand, the proposed multigrid method adds no more than a factor 5/3 to the work in 2D and 58/35 in 3D (ignoring the effort of finding the flow-aligned cell strings). Coding complexity as compared to isotropic semicoarsening is also significantly reduced; in fact, such isotropic semicoarsening seems virtually unfeasible on unstructured grids, while the streamwise semicoarsening is feasible as demonstrated by Mavriplis [7], who selects strings of cells in order to apply line solvers.¹

5. Elliptic/hyperbolic splitting approach

For subsonic flows, the steady Euler equations are known to have a mixed behavior: elliptic for the acoustic components, hyperbolic for the advective components. Conventional multigrid methods based on full coarsening are known to work pre-eminently for isotropic elliptic problems, being capable of achieving $O(N)$ -type convergence, but not for hyperbolic, advection-dominated, problems. In the latter case a remedy is to coarsen the grid only in the advection direction, while retaining the original resolution in the cross-stream direction, i.e., semicoarsening.

Our approach is then to separate the two problem-types at the residual level, and apply the optimal strategy to each residual component: full coarsening to the elliptic part, streamwise semicoarsening to the hyperbolic part (see Fig. 3). Note that this greatly simplifies the grid structure and therefore the multigrid coding compared with the isotropic semicoarsening of Mulder (see Fig. 2).

To decompose the residual, we make use of the van Leer–Lee–Roe (VLR) [17] preconditioning matrix corrected for the effect of the cell-aspect ratio AR [4]; its generic form for grid-aligned subsonic flow, based on the symmetrized Euler equations (with state vector differential $(dp/(\rho a), du, dv, ds)^T$, s denoting entropy), is

$$\mathbf{P}_{\text{VLR}} = \begin{bmatrix} \frac{\tau}{\beta^2} M^2 & -\frac{\tau}{\beta^2} M & 0 & 0 \\ -\frac{\tau}{\beta^2} M & \frac{\tau}{\beta^2} + 1 & 0 & 0 \\ 0 & 0 & \tau & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad (10)$$

where M is the local Mach number, $\beta = \sqrt{1 - M^2}$, and

$$\tau = \frac{\beta}{\beta + AR}. \quad (11)$$

The VLR symmetric preconditioner achieves the lowest possible condition number for all Mach numbers. For 2-D subsonic flow it has the property of decoupling the pair of acoustic equations from the remaining hyperbolic equations; it is unique in this respect [8]. There are other preconditioners that achieve the lowest condition number [3]; among these is a generalization of Turkel's [13] lower-triangular low-speed preconditioner. This one does not separate the elliptic from the hyperbolic equations, but it may be used to

¹ A streamwise line solver definitely is an alternative to a multistage advection scheme.

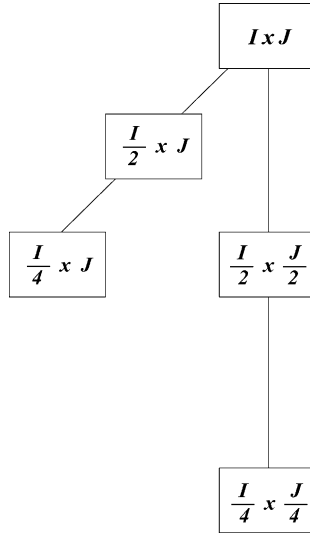


Fig. 3. The grid sequences for the new approach: full/semi-coarsening for the elliptic/hyperbolic residual part.

split the residual into elliptic and hyperbolic contributions. However, there is no practical experience in such use of the Turkel preconditioner, whereas the VLR preconditioner has been routinely used where elliptic/hyperbolic splitting is called for [8,11]; see also [10] which describes the derivation of the VLR preconditioner based on elliptic/hyperbolic decomposition of PDEs. This justifies the choice of the VLR preconditioner for the current application.

Without going into the details of its derivation, we wish to point out that the matrix contains elements and terms in elements proportional to τ , and other elements and terms in elements that do not contain τ . The terms proportional to τ together select the acoustic contributions to the residual, the other terms select the advective contributions. The value of τ is the ratio of the acoustic speed to the flow speed in the preconditioned world. When applying the preconditioner to the Euler equations, setting τ equal to β yields the lowest possible condition number, $K = 1/\beta = 1/\sqrt{1 - M^2}$, down from $K = (M + 1)/\min(M, 1 - M)$ before preconditioning. When applying the preconditioner to an Euler *discretization*, the value $\beta/(\beta + AR)$ is better [3], as it takes into account the effect of artificial dissipation.

The elliptic/hyperbolic decomposition therefore is given by

$$\mathbf{P}_{\text{VLR}} = \mathbf{P}_{\text{VLR}}^{\text{ell}} + \mathbf{P}_{\text{VLR}}^{\text{hyp}}, \tag{12}$$

where

$$\mathbf{P}_{\text{VLR}}^{\text{ell}} = \begin{bmatrix} \frac{\tau}{\beta^2} M^2 & -\frac{\tau}{\beta^2} M & 0 & 0 \\ -\frac{\tau}{\beta^2} M & \frac{\tau}{\beta^2} & 0 & 0 \\ 0 & 0 & \tau & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \quad \mathbf{P}_{\text{VLR}}^{\text{hyp}} = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}. \tag{13}$$

Applying $\mathbf{P}_{\text{VLR}}^{\text{ell}}$ and $\mathbf{P}_{\text{VLR}}^{\text{hyp}}$ to the residual gives the elliptic and hyperbolic part of the residual within each cell, respectively, i.e.,

$$\mathbf{Res}^{\text{ell}} = \mathbf{P}_{\text{VLR}}^{\text{ell}} \mathbf{Res}, \quad \mathbf{Res}^{\text{hyp}} = \mathbf{P}_{\text{VLR}}^{\text{hyp}} \mathbf{Res}. \tag{14}$$

Note that in the residual computation nothing changes; in particular, the flux function (6) is still computed using the *full* preconditioning matrix \mathbf{P}_{VLR} .

Using the above decomposition we are able to compute acoustic and advective changes to the solution separately, and on separate grid sequences. Starting on the finest grid, we update with purely acoustic contributions by using the baseline scheme (8) in which \mathbf{P} is replaced by $\mathbf{P}_{\text{VLR}}^{\text{ell}}$, i.e.,

$$\bar{\mathbf{U}}^{(k+1)} = \bar{\mathbf{U}}^n + \frac{\beta^{(k+1)} \Delta t}{S} \mathbf{P}_{\text{VLR}}^{\text{ell}(k)} \mathbf{Res}^{(k)}. \quad (15)$$

We continue doing so on the fully coarsened grid sequence (cf. Fig. 3), according to the chosen multigrid cycle. Similarly, we update with purely advective contributions by using the baseline scheme (8) in which \mathbf{P} is replaced by $\mathbf{P}_{\text{VLR}}^{\text{hyp}}$, i.e.,

$$\bar{\mathbf{U}}^{(k+1)} = \bar{\mathbf{U}}^n + \frac{\beta^{(k+1)} \Delta t}{S} \mathbf{P}_{\text{VLR}}^{\text{hyp}(k)} \mathbf{Res}^{(k)}, \quad (16)$$

and continue doing so on the streamwise coarsened grid sequence (cf. Fig. 3), according to the chosen multigrid cycle.

For the simple test problem considered in this study we used the simplest possible restriction and prolongation strategies: restriction by conservation over four/two cells for the elliptic/hyperbolic part, and prolongation by bilinear/linear interpolation. The acoustic and advective residuals were treated completely independently throughout the multigrid cycle; only at the finest grid level, where the elliptic and hyperbolic grids are identical, the total solution corrections for the different components were combined into a solution update. The convergence results show that this worked; for more complex problems we would expect the need for tighter coupling between the operators acting on the elliptic and hyperbolic residual components.

6. Numerical results

Our test case is flow over a smooth 5%-high bump in the interval $0 \leq x \leq 1$ described by $y = 0.05 \sin(\pi x)$ (no stagnation points), for which many results are available in the literature [2,6]. The flow is almost perfectly aligned with the grid and conventional multigrid methods fail to achieve $O(N)$ -type convergence.

The domain is defined by $-1 \leq x \leq 3.5$ and $0 \leq y \leq 2$; the boundary at $y = 0$ is treated as solid and the other ones as open. At the boundaries, a ghost-cell approach is used: reflection of the normal velocity while keeping all other state variables the same across the solid wall; prescribing freestream values across open boundaries. The boundary fluxes are then determined by the same Riemann solver as used in the interior. The grid is a structured quadrilateral grid with modest compression toward the wall. A sample grid and solution are shown in Figs. 4 and 5.

A numerical solution is taken to be converged when all L_1 norms of the components of the residual, viz. mass, momentum and energy, are reduced five orders of magnitude from their initial values.

For the multigrid relaxation an $F(1,1)$ -cycle is chosen with the coarsest grid of 8×4 cells; note that the streamwise semicoarsening (essentially in the x -direction) keeps the number of cells in the y -direction the same for all grids. A work unit is defined as the amount of work required to evaluate the full residual on the finest grid. The total number of work units presented in the following results includes four initial relaxations on the finest grid prior to the multigrid cycles.

The convergence histories for various grids (64×32 , 128×64 , 256×128 , 512×256) are shown in Figs. 6 and 7, comparing the results of full-coarsening F -cycle multigrid with a preconditioned four-stage Jacobi relaxation and the new elliptic/hyperbolic decomposition method with the same F -cycle and four-stage Jacobi relaxation for each subproblem. It is evident that the convergence is independent of the grid size in the new method, whereas the conventional method requires an increasing amount of work to converge as the grid gets finer.

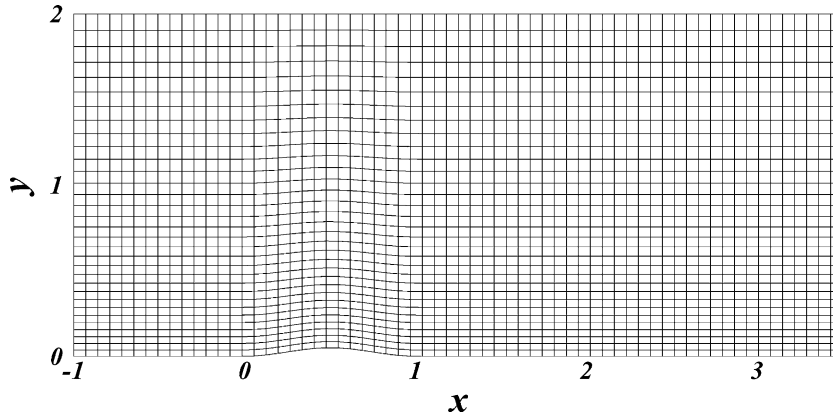


Fig. 4. A 64×32 grid.

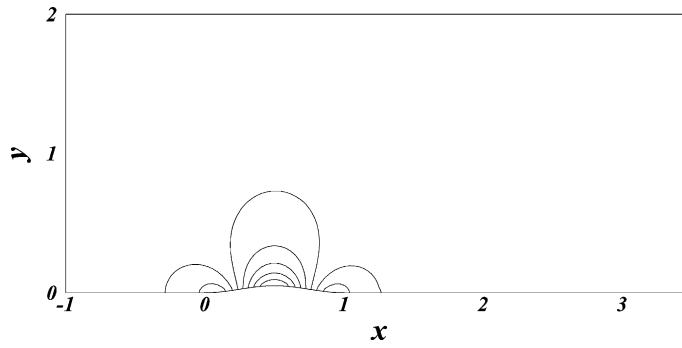


Fig. 5. Mach contours on the 128×64 grid; $M_\infty = 0.1$.

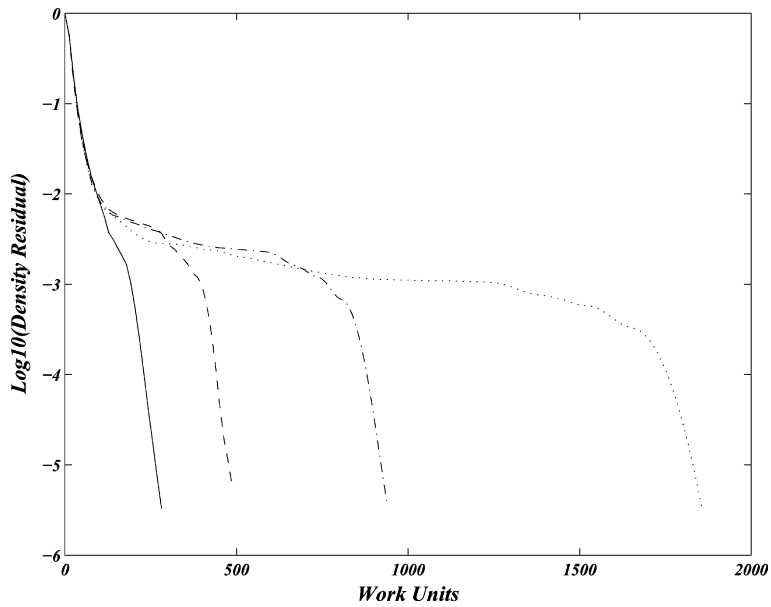


Fig. 6. Convergence histories for full coarsening ($M_\infty = 0.1$). Solid: 64×32 ; dash: 128×64 ; dash-dot: 256×128 ; dot: 512×256 .

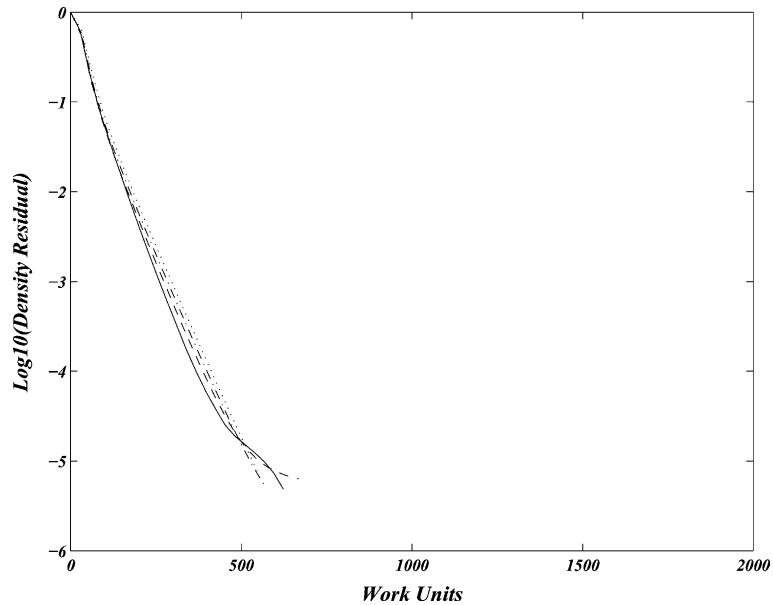


Fig. 7. Convergence histories for the new decomposition approach ($M_\infty = 0.1$). Solid: 64×32 ; dash: 128×64 ; dash-dot: 256×128 ; dot: 512×256 .

In Figs. 8 and 9 convergence histories for the two methods are plotted for different Mach numbers. For both methods the plots coincide, showing independence of Mach number. This is a well-known property of the VLR preconditioner. Without preconditioning, this is not possible, as reported by many researchers [2].

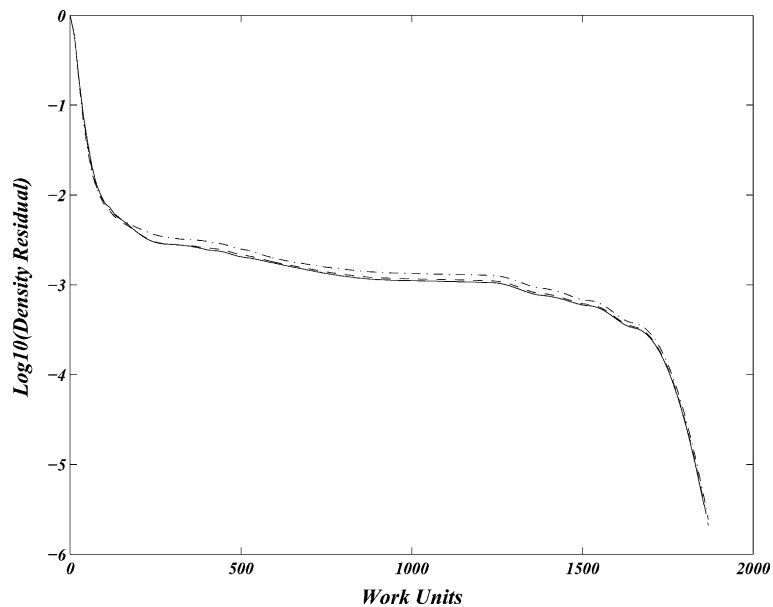


Fig. 8. Convergence histories for full coarsening on 512×256 grid for different Mach numbers. Solid: $M_\infty = 0.1$; dash: $M_\infty = 0.3$; dash-dot: $M_\infty = 0.5$.

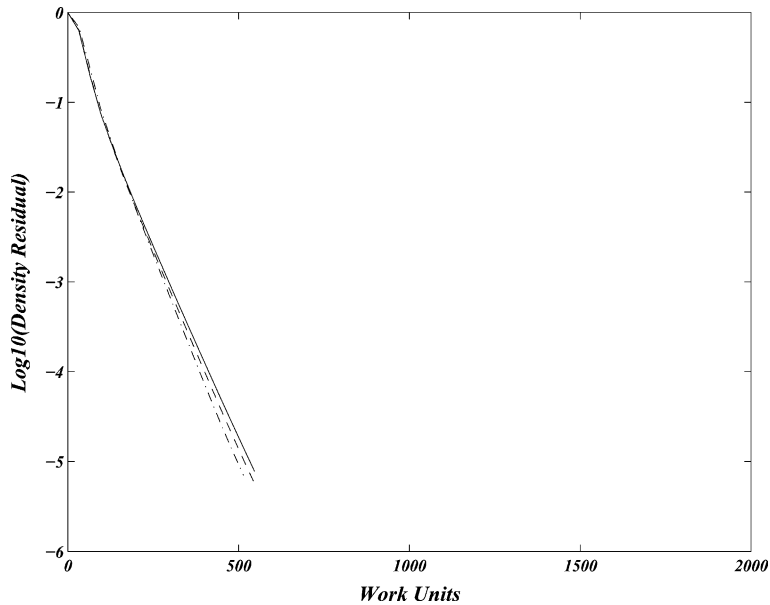


Fig. 9. Convergence histories for the new approach on 512×256 grid for different Mach numbers. Solid: $M_\infty = 0.1$; dash: $M_\infty = 0.3$; dash-dot: $M_\infty = 0.5$.

Lastly, actual CPU times for the two methods are compared in Fig. 10, measured when the programs are run on a single Pentium III (500 MHz) processor. The curves were determined by least-squares linear fits in log–log space. The plots show that the convergence rate is $O(N)$ for the new method, while nearly $O(N^{1.5})$ for full coarsening. Thus, for sufficiently fine grids the new method outperforms the standard method.

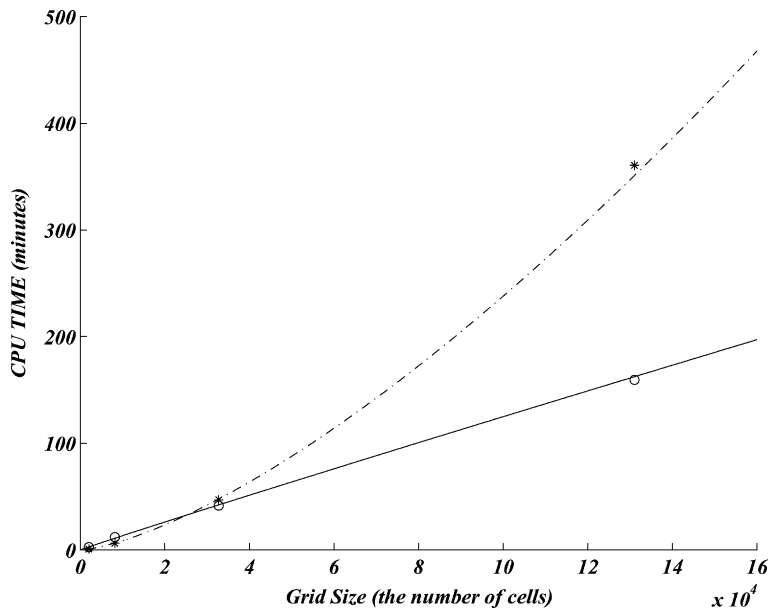


Fig. 10. CPU time versus grid size for $M_\infty = 0.1$. Solid: new approach $CPU = 1.7 \times 10^{-3} \times N^{0.971}$; dash: full coarsening $CPU = 1.53 \times 10^{-5} \times N^{1.44}$ where N is the number of cells.

7. Conclusions and future work

The numerical results presented provide the proof of concept for the current elliptic/hyperbolic multigrid method. Putting it differently: had the proposed strategy not worked for this problem, we might as well forget about making it work for more complex problems. The great advantage of the adopted test problem was that it was a priori known which strings of cells had to be semicoarsened. Thus, the investment in coding streamwise coarsening was minimal.

Obviously, we are just at the start of the development of a more robust method. The following points need to be attended in the future:

- (1) Development of a tightly coupled restriction-prolongation strategy for the elliptic and hyperbolic residual components.
- (2) Formulation of the method on unstructured grids, where one has to determine “on the go” which cells are strongly coupled in the flow direction. Here we can learn from previous experience by Mavriplis [7], who selects strings of cells in order to apply line solvers.
- (3) Overcoming the problem of reduced acoustic-residual smoothing near a sonic line or for large cell-aspect ratios (these problems do not arise with isotropic semicoarsening).
- (4) Blending the method with a semicoarsening strategy for supersonic flow, which by itself still needs development.
- (5) Extending the method to three dimensions. As mentioned above, one of the hyperbolic components cannot be decoupled from the acoustic residual with conventional preconditioning; decoupling is only achieved if the preconditioning operator include differentiation entries. For certain discretizations (multi-D residual-distribution schemes) it may actually be possible to achieve this numerically without leaving the conservation form of the equations. Meanwhile, it remains to be seen how seriously multigrid convergence would be effected if the decomposition were not complete.

Acknowledgements

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References

- [1] W.K. Anderson, J.L. Thomas, B. van Leer, Comparison of finite volume flux vector splittings for the Euler equations, *AIAA J.* 24 (1985) 1453–1460.
- [2] D.L. Darmofal, K. Siu, A robust multigrid algorithm for the Euler equations with local preconditioning and semi-coarsening, *J. Comput. Phys.* 151 (1999) 728–756.
- [3] D. Lee, Design criteria for local Euler preconditioning, *J. Comput. Phys.* 144 (1998) 423–459.
- [4] J.F. Lynn, Multigrid solution of the Euler equations with local preconditioning, Ph.D. Thesis, University of Michigan, 1995.
- [5] J.F. Lynn, B. van Leer, Multi-stage schemes for the Euler and Navier–Stokes equations with optimal smoothing, *AIAA Paper* 93-3355-CP, 1993.
- [6] J.F. Lynn, B. van Leer, D. Lee, Multigrid solution of the Euler equations with local preconditioning, in: *Fifteenth International Conference on Numerical Methods in Fluid Mechanics*, Springer, Berlin, 1996.
- [7] D.J. Mavriplis, On convergence acceleration techniques for unstructured meshes, *AIAA Paper* 98-2966, 1998.
- [8] L.M. Mesaros, Multi-dimensional fluctuation splitting schemes for the Euler equations on unstructured grids, Ph.D. Thesis, University of Michigan, 1995.
- [9] W.A. Mulder, A new approach to convection problems, *J. Comput. Phys.* 83 (1989) 303–323.
- [10] P.L. Roe, H. Nishikawa, K. Kabin, Toward a general theory of local preconditioning, in: *32nd AIAA Fluid Dynamics Conference*, June 24–27, St. Louis, MO, 2002, *AIAA* 02-2956.